

Effective Quantum Plane from Quantum Physics

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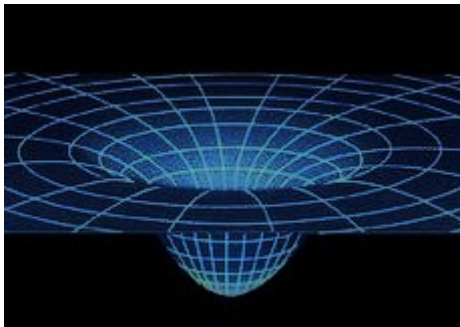
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31st LQP Workshop, Leipzig

- 1 Introduction
- 2 Deformation of QM
- 3 Deformation in QFT
 - Deforming QF with special conformal operator
- 4 Conclusion and Outlook

Motivation for physics on quantized spacetime

- Four known fundamental interactions. Standard model unifies three interactions. Attempts to unify gravity with the SM failed \Rightarrow Quantum Gravity.
- Principles of QM+GRT: classical picture of spacetime breaks down near distances of the Planck length $\lambda_{Pl} = \left(\frac{G\hbar}{c^3}\right)^{1/2}$.

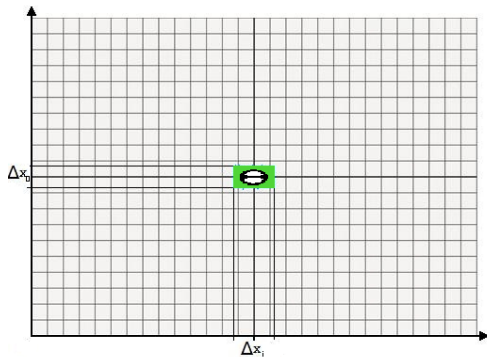


Replace coordinates x^μ by self-adjoint operators \hat{x}^μ obeying

$$[\hat{x}^\mu, \hat{x}^\nu] = i\theta^{\mu\nu}$$

\Rightarrow Uncertainty relations:

$$\Delta x^0 \Delta x^i \geq \theta/2$$



- Obtain quantum plane from QM and QFT

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- Find interpretation from theory for the deformation constant

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X_μ selfadjoint and abelian $[X_\mu, X_\nu] = 0$,

$\implies \exists$ strongly continuous unitary group $V(k) := e^{ik^\mu X_\mu}$

Buchholz, Lechner, Summers '10:

Definition of deformation

Let θ be a skew-symmetric matrix, then the warped convolution $A_{\theta, X}$ of $A \in C^\infty$ is

$$A_{\theta, X} \Phi := (2\pi)^{-d} \int \int dy dk e^{-iyk} V(\theta y)(A)V(-\theta y)V(k)\Phi, \quad \Phi \in \mathcal{D} \subset \mathcal{H}$$

Free Hamiltonian:

$$H_0 = -P_j P^j / (2m)$$

Deformed Hamiltonian with warped convolutions

$$H_{\theta, X} \Psi = -\frac{1}{2m} (P_j + \theta_{jk} X^k) (P^j + \theta^{jr} X_r) \Psi = -\frac{1}{2m} P_j^{\theta, X} P_{\theta, X}^j \Psi$$

The deformed Energy Eigenvalues

Quantized energy values for fixed p_1 !

$$E_{\theta, n} = \frac{p_1^2}{2m} + \left(n + \frac{1}{2} \right) \frac{\theta}{m}, \quad p_1 \in \mathbb{R}, n \in \mathbb{N}.$$

Landau levels

By setting $\theta = eB$ for deformed free Hamiltonian

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Zeemaneffect

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Aharanov Bohm Effect

H_0 deformed with $F_j(X)$ and $\theta = -e\phi_M/2\pi$ induces AB gauge field

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⇒ Obtain any gauge field!

Magnetic Translation Group

Generators of MTG

$$W(\mathbf{a}) = e^{ia^j \pi_j}$$

with canonical momentum operator $\pi_i = P_i - eA_i$ satisfying the CR

$$[\pi_i, \pi_j] = -iF_{ij}.$$

Lemma

For $\theta_{ij} = -\frac{1}{2}F_{ij}$, the deformed generator of the Heisenberg-Weyl group $P_{\theta, X}^j$ is equal to the canonical momentum operator of the MTG π_i .

In lowest LL motion described by $Q_i = X_i + (B^{-1})_{ik} P^k$,

$$[Q_i, Q_j] = 2i(B^{-1})_{ij}.$$

$\implies Q_i$ spans quantum plane $\mathbb{R}_{2B^{-1}}^3$.

Lemma

$X_i^{\theta,P}$ satisfies the commutation relations of the Moyal-Weyl plane $\mathbb{R}_{-2\theta}^3$,

$$[X_i^{\theta,P}, X_j^{\theta,P}] = -2i\theta_{ij}.$$

If $-\theta_{ij}$ is $(B^{-1})_{ij}$, then $X_i^{\theta,P}$ are equal to the guiding center coordinates Q_i .

\implies Idea of lemma can be used in QFT!

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ϕ obeys Klein Gordon Equation

$$\left(\frac{\partial}{\partial x^\mu} \frac{\partial}{\partial x_\mu} + m^2 \right) \phi_0(x) = 0,$$

Solution:

$$\phi_0(x) = \int \left(e^{-ipx} a(\mathbf{p}) + e^{ipx} a^*(\mathbf{p}) \right) d^3\mathbf{p} / 2\omega_{\mathbf{p}}$$

where $\omega_{\mathbf{p}} = \sqrt{-p^i p_i + m^2}$.

Quantum plane from warped convolutions

Doplicher, Fredenhagen and Roberts '01

Realized scalar QF on NC Minkowski spacetime, represented on $\mathcal{V} \otimes \mathcal{H}$.

Grosse and Lechner '07

Represented NC counterpart of scalar QF, on \mathcal{H} instead of $\mathcal{V} \otimes \mathcal{H}$.

Buchholz, Summers '08

Introduced warped convolutions and obtained GL QF by deformation with P_μ .

⇒ **Idea** : Deform QF with FS operator to obtain quantum plane.

Commutation relations

$$\begin{aligned}[P_\rho, K_\mu] &= 2i(\eta_{\rho\mu}D - M_{\rho\mu}), & [K_\rho, M_{\mu\nu}] &= i(\eta_{\rho\mu}K_\nu - \eta_{\rho\nu}K_\mu), \\ [D, K_\mu] &= iK_\mu, & [K_\mu, K_\nu] &= 0\end{aligned}$$

Swieca, Völkel '76: Construct U_R in \mathcal{H}_1

$$K_\mu := U_R P_\mu U_R$$

$\Rightarrow \exists$ str. cont. unitary group $U(b) := e^{ib^\mu K_\mu}$

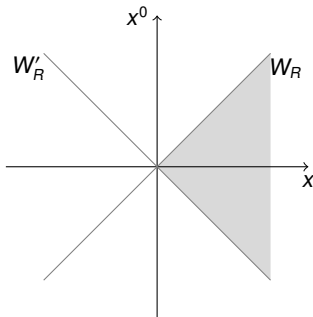
$$\phi_{\theta,K}(f)\Psi := (2\pi)^{-d} \iint dy dk e^{-iyk} U(\theta y)\phi(f)U(-\theta y + k)\Psi, \quad \Psi \in \mathcal{D}(K)$$

- Deformed field fulfills the same bounds as undeformed field
- $\phi_{\theta,K}(f)$ satisfies Reeh-Schlieder theorem but not covariance and locality

Covariance and locality replaced by modified versions!

* AM, J. Math. Phys. 53, 082303 (2012)

Deformed fields defined as QF's on wedge by using map $Q : W \mapsto Q(W)$



Lemma

The special conformal transformations $U_{\theta v}$, with $v \in \mathfrak{spU}$ and θ being admissible, map the right wedge into the right wedge $U_{\theta v}(W_1) \subset W_1$.

Theorem

The family of fields $\phi = \{\phi_W : W \in \mathcal{W}_0\}$ defined by the deformation with K_μ is wedge-covariant, w.r.t. Lorentz group. Furthermore, for $n = 2l + 1$, where $l \in \mathbb{N}_0$, the field ϕ is a wedge-local field on \mathcal{H} .

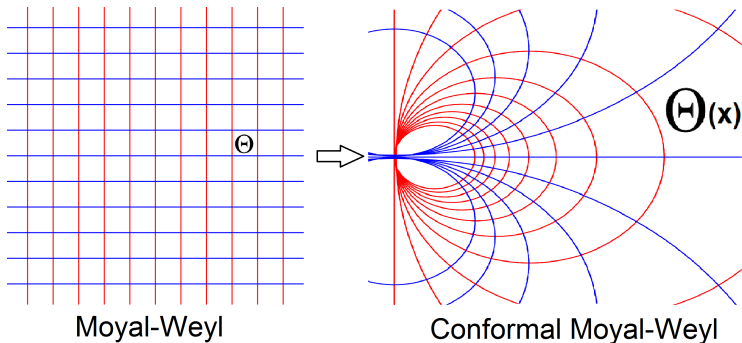
Lemma

Let the deformed product be defined by K_μ . Then the deformed commutator up to third order in θ is

$$[x_\mu \overset{\times_\theta}{,} x_\nu] = -2i\theta_{\mu\nu}x^4 - 4i((\theta x)_\mu x_\nu - (\theta x)_\nu x_\mu)x^2 \equiv i\Theta_{\mu\nu}(x).$$

⇒ Nonconstant noncommutative spacetime.

$$[x_\mu, x_\nu] = -2i\theta_{\mu\nu}x^4 - 4i((\theta x)_\mu x_\nu - (\theta x)_\nu x_\mu) x^2 \equiv i\Theta_{\mu\nu}(x).$$



Deformation of scalar field with operator R_μ

$$R_\mu := (D, M_{01})$$

Proposition

The family of fields $\phi = \{\phi_W : W \in \mathcal{W}_0\}$ defined by the deformation with R_μ is wedge-local, but **not** wedge-covariant.

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- From QM obtained physical effects and physical quantum plane by deformation
- Deformation with SCOP Wightman field with wedge-covariant and wedge-local properties
- Quantum plane from SCOP deformation is nonconstant

- Obtain Stark effect from deformation
- Scattering for deformed QFT's
- Deformation for Fermions and Gauge fields
- Generalize deformation to non Abelian groups
- ...

